

**The two-level local projection
stabilization
as an enriched one-level approach.
A one-dimensional study ^{*}**

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Abstract

The two-level local projection stabilization is considered as a one-level approach in which the enrichments on each element are piecewise polynomial functions. The dimension of the enrichment space can be significantly reduced without losing the convergence order. For example, using continuous piecewise polynomials of degree $r \geq 1$, only one function per cell is needed as enrichment instead of r in the two-level approach. Moreover, in the constant coefficient case, we derive formulas for the user-chosen stabilization parameter which guarantee that the linear part of the solution becomes nodal exact.

^{*}This paper is dedicated to G.I. Shishkin on the occasion of his 70th birthday

1 Introduction

It is well-known that standard Galerkin finite element discretizations applied to convection-diffusion problems show spurious oscillations unless the mesh is adapted to the boundary layers of the solutions [21]. But even in the case of layer adapted meshes it makes sense to use stabilized finite element schemes in order to reduce sensitivities of the solutions on the choice of mesh parameters. Residual based stabilization methods like the streamline upwind Petrov-Galerkin (SUPG) stabilization, proposed in [5] and at first analyzed for a scalar convection-diffusion equation in [19], is a prominent example of stabilized schemes. They rely on adding weighted residuals to the standard Galerkin method to enhance stability without losing consistency.

Recently, local projection stabilization (LPS) [2, 3, 9, 10, 12, 13, 17, 18, 20] methods have become quite popular, in particular because of their commutative properties in optimization problems [4] and stabilization properties similar to those of the SUPG method [11]. In contrast, to residual based stabilization methods the LPS is no longer consistent. However, taking rich enough projection spaces any desired consistency order can be achieved. As shown in [17], the key issue in analyzing the error of LPS schemes is the existence of an interpolation for which the error is orthogonal to the projection space. It turns out, that a local inf-sup condition for the approximation and projection space is sufficient to modify an interpolation into the approximation space in such a way that the additional orthogonality property holds [17]. Two main approaches of LPS have been considered in the literature to fulfil the local inf-sup condition. In the one-level approach, a standard finite element space is chosen as the projection space to guarantee the consistency order. Then, the approximation space is (if necessary) enriched such that the local inf-sup condition holds. In the two-level approach, a standard finite element space is chosen as the approximation space and the projection space is thinned out to a space on the next coarser mesh level to satisfy the local inf-sup condition.

The main objective of this paper is to show that the two-level variant of the LPS can be also considered as an enriched one-level method. This enables us to reduce the degrees of freedom in the two-level method without losing the convergence order. Although this observation is true for any space dimension, we restrict our attention here to the one-dimensional case which gives us the opportunity to study this reduced two-level approach in detail. In particular, we show that in the constant coefficient case the stabilization

parameter can be chosen such that the piecewise linear part of the LPS becomes nodal exact.

In the following, we use the standard notations for Sobolev spaces $H^k(D)$, $H_0^k(D)$, $L^2(D) = H^0(D)$ together with their norms and semi-norms $\|\cdot\|_{k,D}$, $|\cdot|_{k,D}$, and $\|\cdot\|_{0,D}$. We will drop D when $D = (0, 1)$. Throughout this paper C denotes a generic positive constant that is independent of the mesh size.

2 Two Variants of Local Projection Stabilization

We consider the two-point boundary value problem

$$-\varepsilon u'' + bu' + cu = f \quad \text{in } (0, 1), \quad u(0) = u(1) = 0, \quad (2.1)$$

under the assumption

$$c - \frac{1}{2}b' \geq \gamma > 0, \quad (2.2)$$

which guarantees a unique weak solution $u \in H_0^1(0, 1)$. Note that in the interesting case $0 < \varepsilon \ll 1$, the solution exhibits boundary and interior layers whose positions depend on the convection field b .

Let $0 = x_0 < x_1 < \dots < x_N = 1$ be a decomposition \mathcal{M}_h of $[0, 1]$ into macro cells $M \in \mathcal{M}_h$ and h_M the diameter of $M \in \mathcal{M}_h$. In the one-level approach we set $\mathcal{T}_h = \mathcal{M}_h$, i.e. we do not distinguish between a macro cell $M \in \mathcal{M}_h$ and a cell $K \in \mathcal{T}_h$. In the two-level approach, each macro cell $M = [x_i, x_{i+1}]$ is subdivided into two son-cells $K_- = [x_i, x_{i+1/2}]$ and $K_+ = [x_{i+1/2}, x_{i+1}]$ each of diameter $h_K = h_M/2$. Then, all son-cells build the decomposition \mathcal{T}_h . Let $V_h \subset H_0^1(0, 1)$ be a finite element space living on \mathcal{T}_h , D_h be a discontinuous projection space associated with the decomposition \mathcal{M}_h , $\pi_h : L^2(0, 1) \rightarrow D_h$ be the L^2 projection, and $\kappa_h := id - \pi_h$ be a fluctuation operator. The stabilized discrete problem is:

Find $u_h \in V_h$ such that for all $v_h \in V_h$

$$\varepsilon(u_h', v_h') + (bu_h' + cu_h, v_h) + \sum_{M \in \mathcal{M}_h} \tau_M(\kappa_h(bu_h'), \kappa_h(bv_h'))_M = (f, v_h).$$

Herein, (\cdot, \cdot) and $(\cdot, \cdot)_M$ denote the inner product in $L^2(0, 1)$ and $L^2(M)$, respectively and τ_M is a user-chosen stabilization parameter. The bilinear

form associated with the left-hand side is coercive with respect to the mesh-dependent norm

$$|||v||| := \left(\varepsilon |v|_1^2 + \gamma \|v\|_0^2 + \sum_{M \in \mathcal{M}_h} \tau_M \|\kappa_h(bv')\|_{0,M}^2 \right)^{1/2}.$$

In contrast to residual based stabilizations, a consistency error appears whose order depends on choosing τ_M and the projection space D_h .

The key idea in the analysis of the LPS lies in the existence of a special interpolant $j_h : H_0^1(0, 1) \rightarrow V_h$ that displays the usual interpolation properties and satisfies in addition the orthogonality property

$$(w - j_h w, q_h) = 0 \quad \forall w \in H_0^1(0, 1), \quad \forall q_h \in D_h.$$

Using the coercivity of the underlying bilinear form, a rich enough projection space D_h , and the properties of the interpolant j_h , we end up with the error estimate

$$|||u - u_h||| \leq C (\varepsilon^{1/2} + h^{1/2}) h^r |u|_{r+1} \quad (2.3)$$

for $\tau_M \sim h_M$ [1, 3, 16, 17]. The existence of an interpolation j_h with additional orthogonality properties is guaranteed by the following result (adjusted to 1d) [17]:

Theorem 2.1. *Let the local inf-sup condition*

$$\inf_{q_h \in D_h(M)} \sup_{v_h \in Y_h(M)} \frac{(v_h, q_h)_M}{\|v_h\|_{0,M} \|q_h\|_{0,M}} \geq \beta_1 > 0, \quad \forall M \in \mathcal{M}_h \quad (2.4)$$

with $Y_h(M) = \{w_h|_M : w_h \in V_h, w_h = 0 \text{ on } (0, 1) \setminus M\}$ and $D_h(M) = \{r_h|_M : r_h \in D_h\}$ be satisfied. Then there is an interpolation $j_h : H_0^1(0, 1) \rightarrow V_h$ with the orthogonality property

$$(w - j_h w, q_h) = 0, \quad \forall q_h \in D_h, \quad \forall w \in H_0^1(0, 1)$$

and the usual interpolation error estimates.

In order to fulfil all assumptions of the convergence analysis, two different requirements for the pair (V_h, D_h) have to be reconciled:

- D_h has to be rich enough to guarantee a certain order of consistency,

- D_h should be small enough w.r.t. V_h to guarantee $j_h u - u \perp D_h$.

From this two different approaches can be derived:

$$\text{one-level } (V_h^+, D_h) \quad \Leftrightarrow \quad \text{two-level } (V_h, D_{2h}).$$

In the one-level approach, a standard finite element space is chosen as the projection space D_h to guarantee the consistency order. Then, the approximation space V_h is (if necessary) enriched to V_h^+ such that the assumptions of Theorem 2.1 are fulfilled. In the two-level approach, a standard finite element space is chosen as the approximation space V_h and the projection space D_h is thinned out to a space D_{2h} on the next coarser mesh level.

3 Two-level LPS as an Enriched One-level LPS

3.1 Discrete Problem

For some $r \in \mathbb{N}$, let the solution and projection spaces of the two-level LPS be defined by

$$\begin{aligned} V_h &:= \{v_h \in H_0^1(0, 1) : v_h|_K \in P_r(K) \quad \forall K \in \mathcal{T}_h\}, \\ D_{2h} &:= \{q_{2h} \in L^2(0, 1) : q_{2h}|_M \in P_{r-1}(M) \quad \forall M \in \mathcal{M}_h\}. \end{aligned}$$

Our stabilized two-level method is:

Find $u_h \in V_h$ such that for all $v_h \in V_h$

$$a(u_h, v_h) + S(u_h, v_h) = (f, v_h) \quad (3.5)$$

where the bilinear form a and the stabilizing term S are given by

$$\begin{aligned} a(u_h, v_h) &= \varepsilon(u_h', v_h') + (bu_h' + cu_h, v_h), \\ S(u_h, v_h) &= \sum_{M \in \mathcal{M}_h} \tau_M (\kappa_{2h}(bu_h'), \kappa_{2h}(bv_h'))_M. \end{aligned}$$

Herein, $\pi_{2h} : L^2(0, 1) \rightarrow D_{2h}$ denotes the L^2 projection and $\kappa_{2h} := id - \pi_{2h}$ the fluctuation operator. Since the pair (V_h, D_{2h}) of spaces satisfies the assumption of Theorem 2.1 we conclude the error estimate

$$|||u - u_h||| \leq C(\varepsilon^{1/2} + h^{1/2})h^r |u|_{r+1} \quad (3.6)$$

for the solution u_h of (3.5) in the mesh-dependent norm

$$|||v||| := \left(\varepsilon |v|_1^2 + \gamma \|v\|_0^2 + \sum_{M \in \mathcal{M}_h} \tau_M \|\kappa_{2h}(bv')\|_{0,M}^2 \right)^{1/2}.$$

As already mentioned, the keypoint in the error analysis is the existence of an interpolation satisfying an additional orthogonality property. We will see that such an interpolation can be already constructed for a subspace of V_h which allows to create a method with less degrees of freedom than the two-level method but with the same convergence rate. We will construct such a subspace and the associated interpolation on the reference macro in the next subsection.

3.2 Splitting of the Approximation Space

Let $\widehat{M} = (-1, +1)$ be the reference macro, $\widehat{K}_- = (-1, 0)$, $\widehat{K}_+ = (0, +1)$, and $F_M : \widehat{M} \rightarrow M$ the affine mapping of \widehat{M} onto $M \in \mathcal{M}_h$. We define the spaces

$$\widehat{P}_{r,h} = \left\{ \hat{v} \in H^1(\widehat{M}) : \hat{v}|_{\widehat{K}_-} \in P_r(\widehat{K}_-), \hat{v}|_{\widehat{K}_+} \in P_r(\widehat{K}_+) \right\}, \quad \widehat{P}_{r,2h} = P_r(\widehat{M}),$$

where $\dim \widehat{P}_{r,h} = 2r + 1$ and $\dim \widehat{P}_{r,2h} = r + 1$. Consider the set of nodal functionals

$$N_i(\hat{v}) = \int_{-1}^{+1} \hat{v}(\xi) L_i(\xi) d\xi, \quad i = 0, 1, \dots, r-1,$$

$$N_r(\hat{v}) = \hat{v}(-1), \quad N_{r+1}(\hat{v}) = \hat{v}(+1)$$

where L_i , $i = 0, 1, \dots$, denote the Legendre polynomials of degree i on $(-1, +1)$ normalized such that $L_i(1) = 1$. The first r nodal functionals guarantee that a local interpolation $\widehat{J} : H^1(\widehat{M}) \rightarrow \widehat{P}_{r,h}$, defined by $N_i(\hat{v} - \widehat{J}\hat{v}) = 0$ for $i = 0, \dots, r+1$, satisfies the orthogonality property

$$(\hat{v} - \widehat{J}\hat{v}, q)_{\widehat{M}} = 0, \quad \forall q \in P_{r-1}(\widehat{M}), \hat{v} \in H^1(\widehat{M}).$$

The last two nodal functionals secure that the interpolation can be extended to a global continuous interpolation $j_h : H_0^1(0, 1) \rightarrow V_h$ with the desired properties. However, because of $\dim \widehat{P}_{r,2h} < r + 2$ we cannot hope to find an interpolation $\widehat{J} : H^1(\widehat{M}) \rightarrow \widehat{P}_{r,2h}$ into the coarse space $\widehat{P}_{r,2h}$ satisfying all

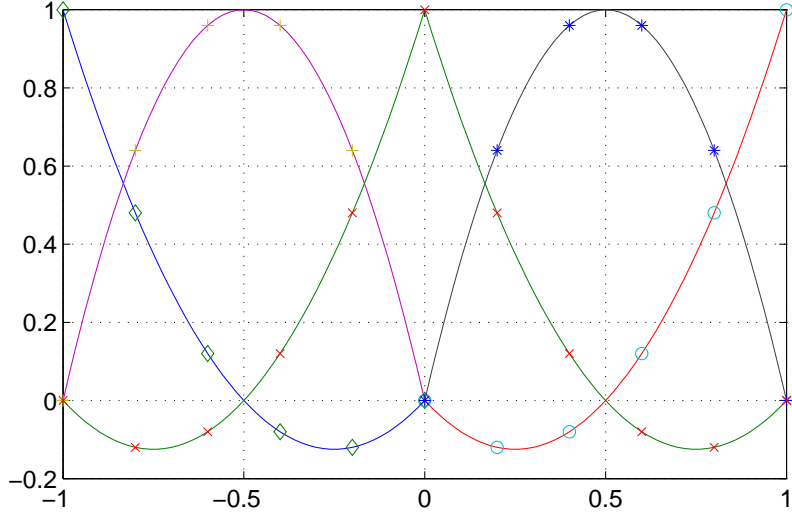


Figure 1: Standard nodal basis of piecewise quadratic functions.

$r + 2$ conditions. We will show that a suitable enrichment of $\widehat{P}_{r,2h}$ by just one additional function is enough to meet these requirements. This additional function is uniquely determined if it is orthogonal (with respect to a certain inner product) to $\widehat{P}_{r,2h}$. Let us consider the following functions

$$\hat{\varphi}_r(x) = \begin{cases} \Lambda_r(x) + \Lambda_{r-1}(x) & x \in [-1, 0] \\ \Lambda_r(-x) + \Lambda_{r-1}(-x) & x \in [0, +1] \end{cases} \quad r \text{ odd,}$$

$$\hat{\varphi}_r(x) = \begin{cases} \Lambda_r(x) - \Lambda_{r-2}(x) & x \in [-1, 0] \\ -(\Lambda_r(-x) - \Lambda_{r-2}(-x)) & x \in [0, +1] \end{cases} \quad r \text{ even,}$$

where Λ_r denotes the Legendre polynomial of degree r on $(-1, 0)$ given by

$$\Lambda_r(x) = L_r(2x + 1) \quad x \in (-1, 0)$$

Furthermore, we introduce the linear mapping $\Phi : \widehat{P}_{r,h} \rightarrow \mathbb{R}^{r+2}$ given by

$$\Phi(\hat{v}) = (N_0(\hat{v}), \dots, N_{r+1}(\hat{v})).$$

Lemma 3.1. *There is the unique splitting*

$$\widehat{P}_{r,h} = \widehat{P}_{r,2h} \oplus \text{span}(\hat{\varphi}_r) \oplus \ker(\Phi). \quad (3.7)$$

The set of nodal functionals N_0, \dots, N_{r+1} is $\widehat{P}_{r,2h}^+$ -unisolvent, where the enriched space is given by $\widehat{P}_{r,2h}^+ = \widehat{P}_{r,2h} \oplus \text{span}(\hat{\varphi}_r)$. Furthermore, the orthogonality property

$$(\hat{\varphi}'_r, \hat{v})_{\widehat{M}} = 0 \quad \text{for all } \hat{v} \in \widehat{P}_{r-1,2h}$$

holds true.

Proof. From the definition of the linear mapping Φ we have

$$\ker(\Phi) = \{\hat{v} \in \widehat{P}_{r,h} : N_i(\hat{v}) = 0, i = 0, \dots, r+1\}$$

which implies $\dim \ker(\Phi) \geq (2r+1) - (r+2) = r-1$.

Now, $\widehat{P}_{r,2h} \cap \ker(\Phi) = \emptyset$, since any function $\hat{v} \in \widehat{P}_{r,2h}$ can be represented as

$$\hat{v} = \sum_{i=0}^r \alpha_i L_i,$$

the orthogonality of the Legendre polynomials yields $\alpha_0 = \dots = \alpha_{r-1} = 0$, and finally we get $0 = N_{r+1}(\hat{v}) = \alpha_r$. Thus, we have $\dim \ker(\Phi) \leq r$.

We prove that $\dim \ker(\Phi) = r-1$ by showing that $\hat{\varphi}_r \notin \widehat{P}_{r,2h}$ and $\widehat{P}_{r,2h}^+ \cap \ker \Phi = \emptyset$. Then, the unique splitting follows immediately. From Rodriguez' formula [6] we get the expansion with respect to powers of x

$$\Lambda_r(x) = L_r(2x+1) = \frac{(2r)!}{2^r(r!)^2} (2x+1)^r + \dots = \frac{(2r)!}{(r!)^2} x^r + \dots,$$

showing that $\hat{\varphi}_r^{(r)}(-0) \neq \hat{\varphi}_r^{(r)}(+0)$ and thus $\hat{\varphi}_r$ cannot be a polynomial on $[-1, +1]$, i.e. $\hat{\varphi}_r \notin \widehat{P}_{r,2h}$. Now we show that $\widehat{P}_{r,2h}^+ \cap \ker \Phi = \emptyset$. This is equivalent to the fact that the set of nodal functionals is $\widehat{P}_{r,2h}^+$ -unisolvent. Since $\dim \widehat{P}_{r,2h}^+ = r+2$ and $r+2$ nodal functionals are given, it suffices to show that

$$\hat{v} \in \widehat{P}_{r,2h}^+, \quad N_i(\hat{v}) = 0, \quad i = 0, \dots, r+1, \quad \Rightarrow \quad \hat{v} = 0.$$

We represent an arbitrary element of $\widehat{P}_{r,2h}^+$ as

$$\hat{v} = \sum_{i=0}^r \alpha_i L_i + \beta \hat{\varphi}_r.$$

Consider the case r odd first. The orthogonality property of Λ_r and Λ_{r-1} on $[-1, 0]$ and of L_k on $[-1, +1]$ imply for $k = 0, 1, \dots, r-2$

$$0 = N_k(\hat{v}) = \alpha_k + \beta \left(\int_{-1}^0 \hat{\varphi}_r(x) L_k(x) dx + \int_0^{+1} \hat{\varphi}_r(x) L_k(x) dx \right) = \alpha_k.$$

Now, from $\hat{\varphi}_r(\pm 1) = 0$ we conclude

$$\begin{aligned} 0 &= N_r(\hat{v}) = (-1)^{r-1}(\alpha_{r-1} - \alpha_r) \\ 0 &= N_{r+1}(\hat{v}) = \alpha_{r-1} + \alpha_r, \end{aligned}$$

consequently $\alpha_{r-1} = \alpha_r = 0$. From the orthogonality properties of Λ_k on $[-1, 0]$ we obtain for r odd

$$\begin{aligned} N_{r-1}(\hat{\varphi}_r) &= \int_{-1}^0 \Lambda_{r-1}(x) L_{r-1}(x) dx + \int_0^{+1} \Lambda_{r-1}(-x) L_{r-1}(x) dx \\ &= \int_{-1}^0 \Lambda_{r-1}(x) L_{r-1}(x) dx + \int_{-1}^0 \Lambda_{r-1}(x) (-1)^{r-1} L_{r-1}(x) dx \\ &= 2 \int_{-1}^0 \Lambda_{r-1}(x) L_{r-1}(x) dx \neq 0. \end{aligned}$$

As $N_{r-1}(\hat{\varphi}_r) \neq 0$ but $N_{r-1}(\hat{v}) = \beta N_{r-1}(\hat{\varphi}_r) = 0$ we get $\beta = 0$.

Consider now the case r even. Again, from orthogonality properties of Λ_r and Λ_{r-1} on $[-1, 0]$ and of L_k on $[-1, +1]$ we get for $k = 0, 1, \dots, r-3$

$$0 = N_k(\hat{v}) = \alpha_k + \beta \left(\int_{-1}^0 \hat{\varphi}_r(x) L_k(x) dx + \int_0^{+1} \hat{\varphi}_r(x) L_k(x) dx \right) = \alpha_k.$$

Orthogonality of Λ_r to L_{r-2} on $[-1, 0]$ yields

$$\begin{aligned} N_{r-2}(\hat{\varphi}_r) &= - \int_{-1}^0 L_{r-2}(x) \Lambda_{r-2}(x) dx + \int_0^{+1} L_{r-2}(x) \Lambda_{r-2}(-x) dx \\ &= - \int_{-1}^0 L_{r-2}(x) \Lambda_{r-2}(x) dx + \int_{-1}^0 (-1)^{r-2} L_{r-2}(x) \Lambda_{r-2}(x) dx = 0. \end{aligned}$$

Thus, $N_{r-2}(\hat{v}) = 0$ implies $\alpha_{r-2} = 0$. From $\hat{\varphi}_r(\pm 1) = 0$ we get

$$\begin{aligned} 0 &= N_r(\hat{v}) = (-1)^{r-1}(\alpha_{r-1} - \alpha_r) \\ 0 &= N_{r+1}(\hat{v}) = \alpha_{r-1} + \alpha_r, \end{aligned}$$

i.e. $\alpha_{r-1} = \alpha_r = 0$. A similar computation as above but now for even r shows

$$\begin{aligned} N_{r-1}(\hat{\varphi}_r) &= - \int_{-1}^0 \Lambda_{r-2}(x) L_{r-1}(x) dx + \int_0^{+1} \Lambda_{r-2}(-x) L_{r-1}(x) dx \\ &= -2 \int_{-1}^0 L_{r-2}(2x+1) L_{r-1}(x) dx = - \int_{-1}^{+1} L_{r-2}(t) L_{r-1}\left(\frac{t-1}{2}\right) dt \\ &= - \frac{(2r-2)!}{2^{r-1}((r-1)!)^2} \int_{-1}^{+1} \left[\left(\frac{t}{2}\right)^{r-1} - \frac{r-1}{2} \left(\frac{t}{2}\right)^{r-2} \right] L_{r-2}(t) dt \neq 0. \end{aligned}$$

Here, we used the orthogonality properties of the Legendre polynomial Λ_r on $[-1, 0]$ and L_{r-2} on $[-1, +1]$, respectively. Now $\beta = 0$ follows from $N_{r-1}(\hat{\varphi}_r) \neq 0$.

It remains to show the orthogonality property. Since $\hat{\varphi}_r \in H_0^1(\widehat{M})$, we obtain by integration by parts

$$(\hat{\varphi}'_r, \hat{v})_{\widehat{M}} = -(\hat{\varphi}_r, \hat{v}')_{\widehat{M}}.$$

For r odd, the right hand side vanishes due to $\hat{v}' \in \widehat{P}_{r-2,2h}$ and the orthogonality of Λ_r and Λ_{r-1} to $\widehat{P}_{r-2,2h}$ on $[-1, 0]$, respectively. For even r , we have

$$\begin{aligned} (\hat{\varphi}'_r, \hat{v})_{\widehat{M}} &= -(\hat{\varphi}_r, \hat{v}')_{\widehat{M}} = \int_{-1}^0 \Lambda_{r-2}(x) \hat{v}'(x) dx - \int_0^{+1} \Lambda_{r-2}(-x) \hat{v}'(x) dx \\ &= A \int_{-1}^0 \Lambda_{r-2}(x) x^{r-2} dx + A \int_0^{-1} \Lambda_{r-2}(x) x^{r-2} dx = 0. \end{aligned}$$

where we assumed an expansion in the form $\hat{v}'(x) = Ax^{r-2} + \dots$ □

For the case $r = 2$ we have

$$\hat{v}_2(x) = \begin{cases} 6x(1+x) & \text{if } x \in [-1, 0], \\ 6x(1-x) & \text{if } x \in [0, +1], \end{cases}$$

and the kernel of Φ is represented by the function

$$\hat{w}(x) = \begin{cases} \frac{3}{2}(1+x)(1+3x) & \text{if } x \in [-1, 0], \\ \frac{3}{2}(1-x)(1-3x) & \text{if } x \in [0, +1]. \end{cases}$$

For $\widehat{P}_{2,h}$ we show the standard nodal basis and a basis corresponding to the splitting of Lemma 3.1 in Figures 1 and 2 .

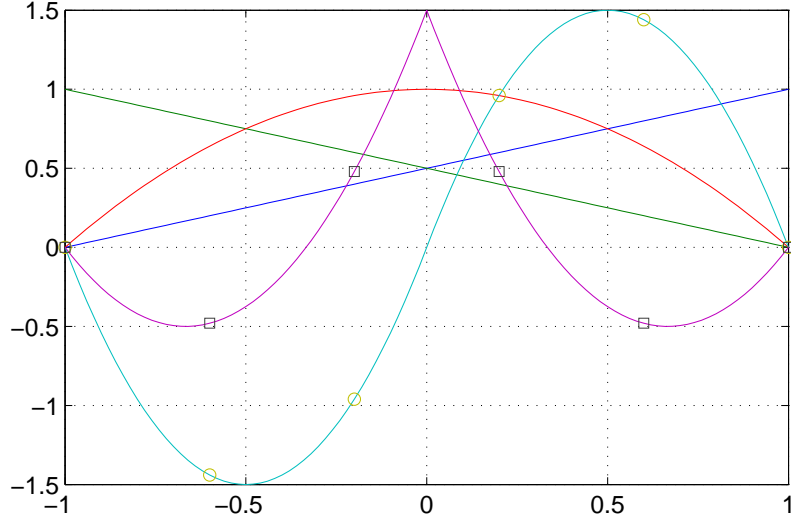


Figure 2: The splitting $\widehat{P}_{2,h} = \widehat{P}_{2,2h} \oplus \text{span}(\widehat{\varphi}_2) \oplus \ker \Phi$. The circles and squares indicate the functions $\widehat{\varphi}_2$ and $\widehat{w} \in \ker \Phi$.

3.3 Reduced Two-level Approach

The results of the previous section motivate us to consider a subspace of V_h for the approximation space combined with the unchanged projection space D_{2h} . Since there is only one function per macro cell added to the same space on the next coarser mesh we consider this type of reduced two-level approach as a one-level approach on \mathcal{M}_h with enrichments of piecewise polynomial functions. Thus, we define the approximation and projection space by

$$\begin{aligned} V_{2h}^+ &:= \{v_h \in H_0^1(0,1) : v_h|_M \in P_{r,2h}^+(M) \quad \forall M \in \mathcal{M}_h\}, \\ D_{2h} &:= \{q_{2h} \in L^2(0,1) : q_{2h}|_M \in P_{r-1}(M) \quad \forall M \in \mathcal{M}_h\} \end{aligned}$$

where $P_{r,2h}^+(M)$ is just the mapped finite element space $\widehat{P}_{r,2h}^+$ introduced in Lemma 3.1. Our reduced two-level discretization is

$$\begin{aligned} \text{Find } u_{2h}^+ \in V_{2h}^+ \text{ such that for all } v_{2h}^+ \in V_{2h}^+ \\ a(u_{2h}^+, v_{2h}^+) + S(u_{2h}^+, v_{2h}^+) = (f, v_{2h}^+). \end{aligned} \quad (3.8)$$

Essential for the error estimation of a solution u_{2h}^+ of (3.8) is the existence of the special interpolation into the reduced finite element space $V_{2h}^+ \subset V_h$.

Lemma 3.2. *There is an interpolation operator $j_h : H_0^1(0, 1) \rightarrow V_{2h}^+$ such that*

$$(j_h w - w, q_{2h}) = 0 \quad \forall q_{2h} \in D_{2h}, w \in H_0^1(0, 1) \quad (3.9)$$

$$|j_h w - w|_{m, M} \leq C h_M^{\ell+1-m} \|w\|_{\ell+1, M} \quad \forall w \in H^{\ell+1}(M), M \in \mathcal{M}_h \quad (3.10)$$

for $\ell = 0, \dots, r, m = 0, 1$.

Proof. On each $M = (x_i, x_{i+1})$ a local interpolant $j_h^M w \in P_{r, 2h}(M)$ is uniquely defined by the $r + 2$ conditions

$$j_h^M w(x_i) = w(x_i), \quad j_h^M w(x_{i+1}) = w(x_{i+1}), \quad (j_h^M w - w, q)_M = 0, \quad \forall q \in P_{r-1}(M)$$

due to Lemma 3.1. The global interpolant $j_h w$ defined by

$$j_h w|_M = j_h^M(w|_M) \quad \forall M \in \mathcal{M}_h$$

belongs to V_{2h}^+ by construction. Since $j_h^M w = w$ for all $w \in P_{r, 2h}(M)$, we obtain (3.10) by means of the Bramble-Hilbert-Lemma. Using $q_{2h}|_M \in P_{r-1}(M)$, we have $(j_h^M w - w, q_{2h})_M = 0$ for all $M \in \mathcal{M}_h$ from which (3.9) follows by summation. \square

Theorem 3.3. *Let u be the weak solution of (2.1) and u_{2h}^+ the the solution of the reduced two-level method (3.8) for $\tau_M \sim h_M$, respectively. Then, the error estimate*

$$|||u - u_{2h}^+||| \leq C \left(\sum_{M \in \mathcal{M}_h} (\varepsilon + h_M) h_M^{2r} \|u\|_{r+1, M}^2 \right)^{1/2}$$

holds provided that $u \in H_0^1(0, 1) \cap H^{r+1}(0, 1)$.

Proof. The proof is analog to that given in [24, Theorem 2]. \square

3.4 Elimination of Enrichments

In the following, we consider the reduced two-level approach (V_{2h}^+, D_{2h}) in the special case that $b = \text{const}$, $c = 0$, and f piecewise $P_{r-1, 2h}$. As in [24] we want to eliminate the enrichments locally. However, in contrast to the one-level approach, we have to deal with enrichments by piecewise polynomial functions

$$\varphi_{r, M}(x) := \hat{\varphi}_r \left(\frac{2x - x_i - x_{i+1}}{x_{i+1} - x_i} \right), \quad x \in M = [x_i, x_{i+1}],$$

which need more care compared to polynomials.

Clearly, we have $\varphi_{r,M}|_M \in H_0^1(M)$. We split the approximation space into the direct sum

$$V_{2h}^+ = V_{2h} \oplus B_h, \quad B_h = \bigoplus_{M \in \mathcal{M}_h} \text{span } \varphi_{r,M},$$

$$V_{2h} := \{v_{2h} \in H_0^1(0,1) : v_{2h}|_M \in P_r(M) \quad \forall M \in \mathcal{M}_h\}.$$

The direct sum $V_{2h}^+ = V_{2h} \oplus B_h$ generates a unique splitting of the solution $u_{2h}^+ \in V_{2h}^+$ of the local projection stabilization in the same way

$$u_{2h}^+ = u_{2h} + \sum_{M \in \mathcal{M}_h} u_M \varphi_{r,M}.$$

Our aim is to reformulate the reduced two-level scheme in terms of this splitting. For this we consider first some terms appearing in the LPS approach. Taking into consideration that for any $v_{2h} \in V_{2h}$ we have $v'_{2h}|_M \in P_{r-1}(M)$, the L_2 projection becomes $\pi_{2h} v'_{2h} = v'_{2h}$ thus the fluctuation $\kappa_{2h}(v'_{2h})$ vanishes. Furthermore, we obtain

$$(b \varphi'_{r,M}, \varphi_{r,M}) = \frac{1}{2}(b, (\varphi_{r,M}^2)') = -\frac{1}{2}(b', \varphi_{r,M}^2) = 0.$$

Now, the reduced two-level discretization (3.8) can be rewritten as

Find $u_{2h} \in V_{2h}$ and $\{u_M\} \in \mathbb{R}^N$ such that for all $v_{2h} \in V_{2h}$, $M \in \mathcal{M}_h$

$$\varepsilon(u'_{2h}, v'_{2h}) + (bu'_{2h}, v_{2h}) + \sum_{M \in \mathcal{M}_h} u_M (b \varphi'_{r,M}, v_{2h}) = (f, v_{2h}), \quad (3.11)$$

$$u_M [\varepsilon(\varphi'_{r,M}, \varphi'_{r,M}) + \tau_M (\kappa_{2h}(b \varphi'_{r,M}), \kappa_{2h}(b \varphi'_{r,M}))] = (f - bu'_{2h}, \varphi_{r,M}).$$

Using the orthogonality $\varphi'_{r,M} \perp P_{r-1}(M)$ stated in Lemma 3.1, we can see that

$$(\pi_{2h} \varphi'_{r,M}, q)_M = (\varphi'_{r,M}, q)_M = 0 \quad \forall q \in P_{r-1}(M),$$

thus the projection $\pi_{2h} \varphi'_{r,M}$ vanishes and $\kappa_{2h} \varphi'_{r,M} = \varphi'_{r,M}$. Then, the second set of equations of the reduced two-level discretization (3.11) reads

$$u_M [\varepsilon + \tau_M b^2] (\varphi'_{r,M}, \varphi'_{r,M}) = (f - bu'_{2h}, \varphi_{r,M}) \quad \forall M \in \mathcal{M}_h. \quad (3.12)$$

Next we intend to apply following Lemma to the right hand side of (3.12).

Lemma 3.4. *There is a $\psi_M \in H^r(M)$ with $\psi_M|_{K_\pm} \in P_{2r-1}(K_\pm)$ satisfying*

$$\psi_M^{(r-1)} = \varphi_{r,M} \quad \text{in } M, \quad \psi_M^{(j)} = 0 \quad \text{on } \partial M, \quad j = 0, 1, \dots, r-1.$$

Proof. We show that on the reference macro $\widehat{M} = [-1, +1]$ there is a function $\widehat{\psi} \in H^r(\widehat{M})$ with

$$\widehat{\psi}^{(r-1)} = \widehat{\varphi}_r \quad \text{in } \widehat{M}, \quad \widehat{\psi}^{(j)}(\pm 1) = 0, \quad j = 0, \dots, r-1. \quad (3.13)$$

Thanks to the formula of Rodriguez we derive an explicit representation. Consider the case of odd r first. Then, a careful check shows that

$$\widehat{\psi}(x) = \begin{cases} \frac{1}{r!} \frac{d}{dx} [x(1+x)]^r + \frac{1}{(r-1)!} [x(1+x)]^{r-1} & x \in [-1, 0], \\ \frac{1}{r!} \frac{d}{dx} [x(1-x)]^r + \frac{1}{(r-1)!} [x(1-x)]^{r-1} & x \in [0, +1] \end{cases}$$

satisfies all requirements of the Lemma. Similarly, for even r we derive a formula starting with

$$\widehat{\psi}^{(r)}(x) = \widehat{v}'_r(x) = 2(2r-1) \begin{cases} L_{r-1}(1+2x) & x \in [-1, 0], \\ L_{r-1}(1-2x) & x \in [0, +1], \end{cases}$$

integrating this representation $(r-1)$ -times, and using the conditions at $x = -1$. As a result we obtain

$$\widehat{\psi}'(x) = \frac{2(2r-1)}{(r-1)!} \begin{cases} [x(1+x)]^{r-1} & x \in [-1, 0], \\ [x(1-x)]^{r-1} & x \in [0, +1]. \end{cases}$$

Taking into consideration that $\widehat{\psi}'(-t) = -\widehat{\psi}'(t)$ and setting

$$\widehat{\psi}(x) := \int_{-1}^x \widehat{\psi}'(t) dt$$

we have $\widehat{\psi}$ satisfying (3.13).

Using the transformation $F_M : \widehat{M} \rightarrow M$, we see that the function

$$\psi_M(x) = \left(\frac{h_M}{2}\right)^{r-1} \widehat{\psi}\left(\frac{2x - x_i - x_{i+1}}{x_{i+1} - x_i}\right)$$

satisfies all requirements of the Lemma. □

Now, using the fact that $f - bu'_{2h}$ is a piecewise polynomial function of degree less than or equal to $r - 1$ with respect to \mathcal{M}_h , we can integrate by parts

$$(f - bu'_{2h}, \varphi_{r,M}) = (f - bu'_{2h}, \psi_M^{(r-1)}) = (-1)^{r-1} ((f - bu'_{2h})^{(r-1)}, \psi_M)$$

and obtain from (3.12) the representation

$$u_M = (f - bu'_{2h})^{(r-1)}|_M \frac{(-1)^{r-1} (1, \psi_M)}{(\varepsilon + \tau_K b^2) |\varphi_{r,M}|_{1,M}^2}.$$

Using this to eliminate u_M in (3.11) we end up with the following method:

Find $u_{2h} \in V_{2h}$ such that for all $v_{2h} \in V_{2h}$

$$\begin{aligned} \varepsilon(u'_{2h}, v'_{2h}) + (bu'_{2h}, v_{2h}) + \sum_{M \in \mathcal{M}_h} \gamma_M ((bu'_{2h})^{(r-1)}, (bv'_{2h})^{(r-1)})_M \\ = (f, v_{2h}) + \sum_{M \in \mathcal{M}_h} \gamma_M (f^{(r-1)}, (bv'_{2h})^{(r-1)})_M \end{aligned} \quad (3.14)$$

where the parameter γ_M is related to the parameter τ_M of the LPS in the following way

$$\gamma_M = \frac{(1, \psi_M)^2}{(\varepsilon + \tau_M b^2) h_M |\varphi_{r,M}|_{1,M}^2}. \quad (3.15)$$

For the considered case $b = \text{const}$, $c = 0$, and f piecewise polynomial of degree $r - 1$, the method (3.14) is identical with the differentiated residual method (DRM) which has been studied already in [23, 24].

Theorem 3.5. *Assume $b = \text{const}$, $c = 0$, and f piecewise $P_{r-1,2h}$ and let the approximation space $V_{2h} = P_{r,2h} \cap H_0^1(0, 1)$ on each cell $M \in \mathcal{M}_h$ be locally enriched by $\varphi_{r,M}$. Then, eliminating the enrichment in the reduced two-level LPS results in the $P_{r,2h}$ -DRM. The associated stabilization parameter τ_M and γ_M are related by (3.15).*

Remark. It is well-known that by an appropriately chosen stabilization parameter γ_M nodal exactness of the piecewise linear part of the solution can be achieved for the DRM in the case $b = \text{const}$, $c = 0$, $f = \text{const}$. Therefore, the formula (3.15) allows to choose the stabilization parameters τ_M in the

reduced two-level LPS is such a way that the piecewise linear part of the $P_{r,2h}^+$ solution is nodal exact. For this we compute

$$(1, \psi_M) = \begin{cases} \frac{2(r-1)!}{(2r-1)!} \left(\frac{h_M}{2}\right)^r & r \text{ odd} \\ -\frac{(r-2)!}{(2r-3)!} \left(\frac{h_M}{2}\right)^r & r \text{ even} \end{cases}, \quad |\varphi_{r,M}|_{1,M}^2 = \begin{cases} \frac{16r^2}{h_M} & r \text{ odd} \\ \frac{16(2r-1)}{h_M} & r \text{ even} \end{cases}$$

and use the formulas defining $\gamma_{r,M}$ from [24]:

$$\tau_{r,M} = \frac{h_M}{\alpha_r b \Phi_r(q_M)} - \frac{\varepsilon}{b^2}, \quad q_M = \frac{bh_M}{2\varepsilon}, \quad \alpha_r = \begin{cases} 2^{2r+1} r^2 & r \text{ odd}, \\ 2^{2r+1}/(2r-1) & r \text{ even}, \end{cases}$$

$$\Phi_{r+1}(q) = \frac{1}{\Phi_r(q)} - \frac{2r+1}{q}, \quad \Phi_1(q) = \coth q - \frac{1}{q}.$$

4 Numerical Examples

In our first test example we choose $b = 1$, $c = 0$, $f = 1$, and $\varepsilon = 10^{-7}$. Apart from an exponential layer near $x = 1$ the solution can be approximated by $u_{asympt}(x) = x$. The stabilization parameter $\tau_{r,M}$ depends on the polynomial degree $r \geq 1$ and is chosen as in Section 3.4. In Figure 3 we present the results for piecewise quadratic approximations on a uniform macro mesh with $2h = 1/N$, $N = 20$ and $N = 40$. We clearly see the nodal exactness of the linear part (marked by stars) of the computed solution and that oscillations are restricted to the boundary layer region for the reduced two-level approach (left). For the classical two-level variant (right) we have taken also the formulas for $\tau_{r,M}$ given in Section 3.4 but with a modified (experimentally fitted) $\alpha_2 = 13.856$ instead of $\alpha_2 = 32/3$. Note that for this variant nodal exactness cannot be guaranteed. The localization of oscillations on the boundary layer region seems to be less pronounced. For the case of piecewise cubic approximations both variants show a similar localization behaviour, however the amplitudes are much larger for the reduced two-level approach. Next we study the influence of layer adapted meshes. We computed the solution on a uniform mesh of $N = 10$ macro cells and subdivided the last element into two cells, and repeated this approach resulting into eleven and twelve macro cells, respectively. The first three diagrams of Figure 4 show the results for the piecewise cubic reduced two-level method. Oscillations are concentrated

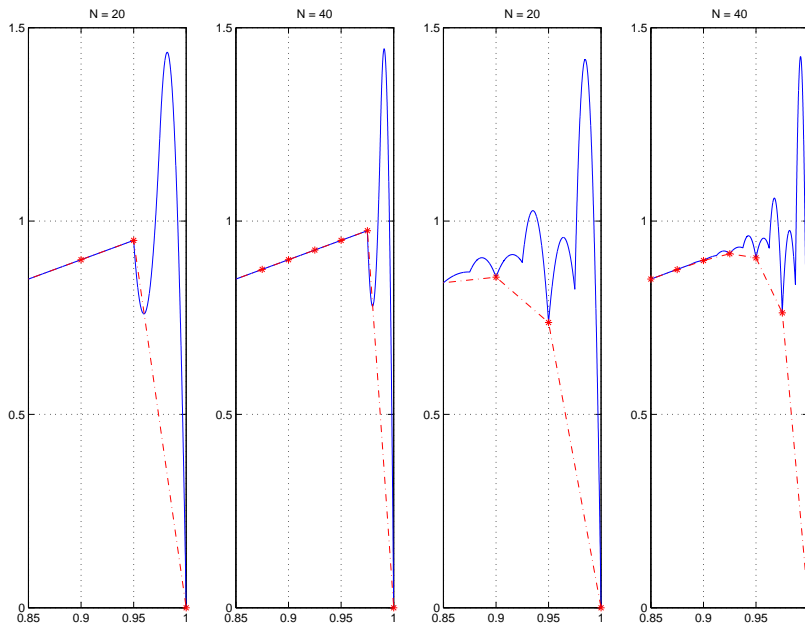


Figure 3: Piecewise quadratic two-level reduced (left) and two-level (right) method. The stars display the linear part of the solutions.

on the last macro cell and their amplitudes are almost constant. The fourth diagram shows the result on a Shishkin mesh with the same number of degrees of freedom $N = 12$. A Shishkin mesh is a piecewise uniform mesh with a transition point at $1 - (3/2)\varepsilon \log N$, see [7, 8, 14, 15, 22]. Thus, the LPS is able to localize oscillations on the layer region but in order to suppress their amplitudes layer adapted meshes seem to be necessary. In the second example we choose $\Omega = (-1, +1)$, $b = -|x|$, $c = 1/2$, $f = 0$, $u(-1) = 1$, $u(1) = 2$. An exponential layer is located near $x = -1$, there is also an inner layer in the first derivative at $x = 0$. In Figure 5 we present the results for the two-level method for a piecewise cubic approximation on a macro mesh of $N = 80$ cells. The main characteristic of LPS, to localize oscillations on layer regions, holds true also in the case of non-constant coefficients. Moreover, no oscillation near the inner layer in the derivative are observed.

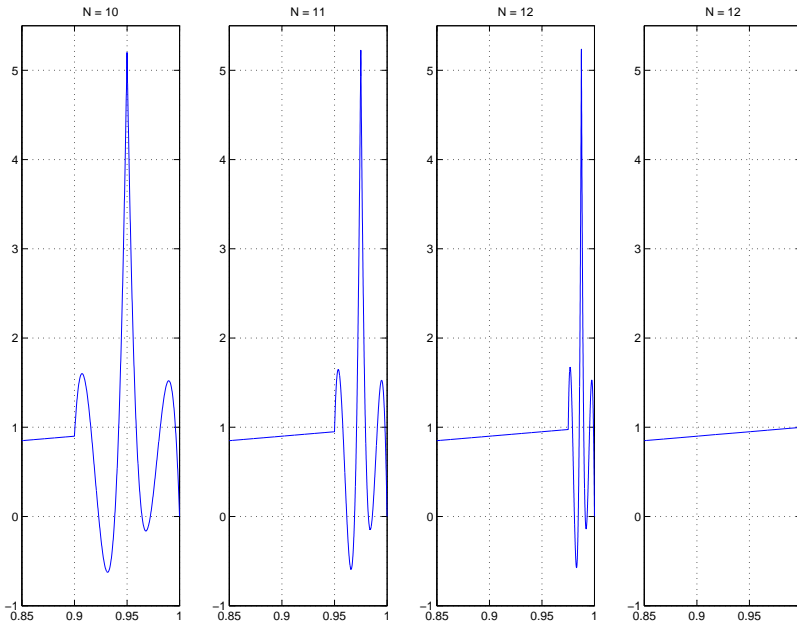


Figure 4: Piecewise cubic reduced two-level method on adapted refined meshes and on a Shishkin mesh (see the right most panel).

5 Conclusions

We have exposed the close relationship of the two variants of the local projection stabilization. In particular, the two-level method can be regarded as a one-level approach on the coarser mesh using piecewise polynomial enrichments. It turns out that from this point of view the two-level approach uses larger enrichments than needed to construct a certain interpolant which guarantees optimal order of convergence. One additional degree of freedom per macro cell is sufficient which corresponds to the one-level approach studied in [24]. Eliminating the enrichments in the constant coefficient case (b, f constant, $c = 0$), we obtain the differentiated residual method [23] for which optimal choices of the stabilization parameters are known [24] leading to the nodal exactness of the linear part of the solution.

Although these investigations are restricted to the one-dimensional case we think they provide some insight into the better understanding of the stabilizing properties of the popular local projection stabilization in the multi-dimensional case.

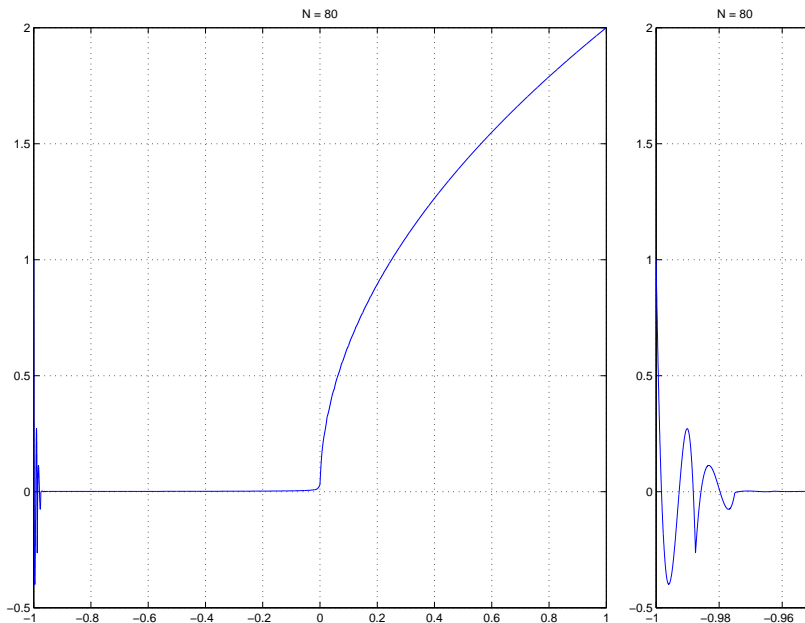


Figure 5: Piecewise cubic two-level method for a non-constant coefficient case and the close-up range of the layer region.

References

- [1] R. Becker and M. Braack. A two-level stabilization scheme for the Navier-Stokes equations. In *Numerical mathematics and advanced applications*, pages 123–130. Springer, Berlin, 2004.
- [2] M. Braack. A stabilized finite element scheme for the Navier-Stokes equations on quadrilateral anisotropic meshes. *Math. Model. Numer. Anal. M2AN*, 42:903–924, 2008.
- [3] M. Braack and E. Burman. Local projection stabilization for the Oseen problem and its interpretation as a variational multiscale method. *SIAM J. Numer. Anal.*, 43:2544–2566, 2006.
- [4] M. Braack and G. Lube. Finite elements with local projection stabilization for incompressible flow problems. *J. Comput. Math.*, 27:116–147, 2009.

- [5] A.N. Brooks and T.J.R. Hughes. Streamline upwind/Petrov-Galerkin formulations for convection dominated flows with particular emphasis on the incompressible Navier-Stokes equations. *Comput. Methods Appl. Mech. Engrg.*, 32:199–259, 1982.
- [6] A. Erdlyi (Ed.). *Higher Transcendental Functions*. vol. II. Based, in part, on notes left by Harry Bateman and compiled by the Staff of the Bateman Manuscript Project. (Repr. of the orig. 1953 publ. by McGraw-Hill Book Company, Inc., New York), Bateman Manuscript Project, California Institute of Technology, Robert E. Krieger Publishing Company, Malabar, Florida, 1981.
- [7] P.A. Farrell, A.F. Hegarty, J.J.H. Miller, E. O’Riordan, G.I. Shishkin. *Robust computational techniques for boundary layers*. Applied Mathematics (Boca Raton), **16**, Chapman & Hall/CRC, Boca Raton, FL, 2000.
- [8] S. Franz and G. Matthies. Local projection stabilization on S-type meshes for convection-diffusion problems with characteristic layers. *Preprint MATH-NM-07-2008, TU Dresden*, 2008.
- [9] J.-L. Guermond. Stabilization of Galerkin approximations of transport equations by subgrid modeling. *M2AN Math. Model. Numer. Anal.*, 33:1293–1316, 1999.
- [10] T.J.R. Hughes and G. Sangalli. Variational multiscale analysis: the fine-scale Green’s function, projection, optimization, localization, and stabilized methods. *SIAM J. Numer. Anal.*, 45(2):539–557, 2007.
- [11] P. Knobloch and L. Tobiska. On the stability of finite element discretizations of convection-diffusion-reaction equations. *IMA J. Numer. Anal.*, (to appear), 2009.
- [12] W.J. Layton. A connection between subgrid scale eddy viscosity and mixed methods. *Applied Mathematics and Computation*, 133:147–157, 2002.
- [13] G. Lube, G. Rapin, and J. Löwe. Local projection stabilization of finite element methods for incompressible flows. In Karl Kunisch, Günther Of, and Olaf Steinbach, editors, *Numerical mathematics and advanced*

- applications. Proceedings of the 7th European Conference (ENUMATH 2007) held in Graz, September 10-14, 2007*, pages 481–488, Berlin, 2008. Springer-Verlag.
- [14] G. Matthies. Local projection methods on layer adapted meshes for higher order discretizations of convection-difusion problems. *Preprint July 1, 2008, Ruhr University Bochum*, 2008.
- [15] G. Matthies. Local projection stabilization for higher order discretizations of convection-difusion problems on Shishkin meshes. *Adv. Comput. Math.*, 2008. doi: 10.1007/s10444-008-9070-y.
- [16] G. Matthies, P. Skrzypacz, and L. Tobiska. Stabilization of local projection type applied to convection-diffusion problems with mixed boundary conditions. *ETNA*, 32: 90–105, 2008.
- [17] G. Matthies, P. Skrzypacz, and L. Tobiska. A unified convergence analysis for local projection stabilisations applied to the Oseen problem. *Math. Model. Numer. Anal. M2AN*, 41(4):713–742, 2007.
- [18] G. Matthies and L. Tobiska. Local projection type stabilisation applied to inf-sup stable discretisations of the Oseen problem. Preprint 07-47, Fakultät für Mathematik, University Magdeburg, 2007.
- [19] U. Nävert. *A finite element method for convection-diffusion problems*. PhD thesis, Chalmers University of Technology, Göteborg, 1982.
- [20] G. Rapin, G. Lube, and J. Löwe. Applying local projection stabilization to inf-sup stable elements. In Karl Kunisch, Günther Of, and Olaf Steinbach, editors, *Numerical mathematics and advanced applications. Proceedings of the 7th European Conference (ENUMATH 2007) held in Graz, September 10-14, 2007*, pages 521–528, Berlin, 2008. Springer-Verlag.
- [21] H.-G. Roos, M. Stynes, and L. Tobiska. *Robust numerical methods for singularly perturbed differential equations. Convection-diffusion-reaction and flow problems*. Number 24 in SCM. Springer-Verlag, 2008.
- [22] G.I. Shishkin and L.P. Shishkina. *Difference methods for singular perturbation problems*. Chapman & Hall/CRC Monographs and Surveys in Pure and Applied Mathematics, **140** CRC Press, 2009.

- [23] L. Tobiska. Analysis of a new stabilized higher order finite element method for advection-diffusion equations. *Comput. Methods Appl. Mech. Engrg.*, 196:538–550, 2006.
- [24] L. Tobiska. On the relationship of local projection stabilization to other stabilized methods for one-dimensional advection-diffusion equations. *Comput. Methods Appl. Mech. Engrg.*, 198:831–837, 2009.